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THE AVERAGE METHOD IS MUCH BETTER THAN AVERAGE

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Dedicated to Professor Barna Szabó on the occasion of his eighty-fifth birthday

Abstract. Operator splitting is a powerful method for the numerical investigation of complex time-dependent models, where the stationary (elliptic) part consists of a sum of several structurally simpler sub-operators. As an alternative to the classical splitting methods, a new splitting scheme is proposed here, the Average Method with sequential splitting. In this method, a decomposition of the original problem is sought in terms of commuting matrices. Wedemonstrate that third-order accuracy can be achieved with the Average Method. The computational performance of the method is investigated, yielding run times 1-2 orders of magnitude faster than traditional methods.

Mathematical Subject Classification: 65L05, 76G25

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1. Introduction

Operator splitting entails the decomposition (splitting) of the spatial differential operator of the problem into a sum of different sub-operators having simpler forms. Operator splitting methods are commonly used in many applications. The decomposition can be motivated mathematically (equations of different types, elliptic, parabolic, etc.) or by the presence of subsystems described by different areas of physics (coupled flow-structure or thermo-mechanical problems, for example).

McLahlan and Quispel [1] survey splitting methods for the numerical integration of ODE's. A nice exposition of splitting methods can be found in [2]. Different communities use different names for the same concept. Operator splitting is also called staggered methods (schemes), decomposition, co-simulation, etc. Gu and Asada [3] discuss the concept of co-simulation, which refers to simultaneous numerical solution

(discretization) of multiple interacting subsystem. A simple mathematical model for co-simulation is proposed in [4], together with the study of the resulting stability charts. Csomós and Nickel introduce splitting methods for delay equations in an abstract setting [5] and prove the convergence of the method.

The structure of this paper is the following. In Section 2 we discuss two basic splitting methods: sequential and Strang-Marchuk splitting. Then we introduce the Average Method in Section 3. In the same section we discuss the possible reduction of the terms needed for the Average Method by using a matrix decomposition of pairwise commuting matrices. In Section 4 we state a condition that makes the Average Method with sequential splitting third-order in accuracy. Here we also show that third-order accuracy cannot be achieved when basing the Average Method on Strang-Marchuk splitting. In Section 5 an example problem is given, motivated by an aerodynamic model. In Section 6 results of eight runtimes of the various methods. Section 7 summarizes the findings of this work.

2. Sequential and Strang-Marchuk splitting

This section is based on [2]. We consider the following Cauchy problem in \mathbb{R}^m

$$\begin{cases} \dot{y}(t) = Ay(t) = \sum_{i=1}^{d} A_i y(t) & t \in [0, T) \\ y(0) = y_0, \end{cases}$$
 (1)

where $y:[0,T]\to\mathbb{R}^m$ is the unknown function, $y_0\in\mathbb{R}^m$ is the given initial vector, $A_i\in\mathbb{R}^{m\times m}$ $(i=1,\ldots,d)$ are matrices.

The exact solution of the Cauchy problem (1) can be written directly as

$$y(t) = \exp(tA)y(0). \tag{2}$$

Our aim is to approximate the exact solution (2) numerically on the grid

$$\omega_h = \left\{ t_n = n \cdot h, h = \frac{T}{N}, n = 0, 1, ..., N \right\}.$$
 (3)

In sequential splitting we decompose the original problem (1) into d sub-problems (i = 1, 2, ..., d)

$$\begin{cases} \dot{y}_{i}^{n}(t) = A_{i}y_{i}^{n}(t), & t \in ((n-1)h, nh], \\ y_{i}^{n}((n-1)h) = y_{i-1}^{n}(nh). \end{cases}$$
(4)

The solution is

$$y_{seq}^N(nh) = y_d^n(nh), (5)$$

where $y_0^n(nh) = y_{seq}^N((n-1)h)$, and $y_{seq}^N(0) = y(0) = y_0$.

Sequential splitting is a first-order method.

The main difference between sequential and *Strang-Marchuk splitting* is that the latter computes the values in the midpoints of the subintervals. First

$$\begin{cases} \dot{y}_{i}^{n}(t) = A_{i} y_{i}^{n}(t), & t \in \left((n-1)h, (n-\frac{1}{2})h\right], \\ y_{i}^{n}\left((n-1)h\right) = y_{i-1}^{n}\left((n-\frac{1}{2})h\right) & \end{cases}$$
(6)

is computed for $i = 1, 2, \dots, d-1$. For i = d the slightly different

$$\begin{cases}
\dot{y}_d^n(t) = A_d y_d^n(t), & t \in ((n-1)h, nh], \\
y_d^n((n-1)h) = y_{d-1}^n((n-\frac{1}{2})h)
\end{cases}$$
(7)

is evaluated. For $i=d+1,\ldots,2d-1$ the following formula is used

$$\begin{cases} \dot{y}_i^n(t) = A_i y_i^n(t), & t \in \left((n - \frac{1}{2})h, nh \right], \\ y_i^n((n - \frac{1}{2})h) = y_{i-1}^n(nh). \end{cases}$$
(8)

The solutions are given by

$$y_{SM}^N(nh) = y_{2d-1}^n(nh),$$

where
$$y_0^n((n-\frac{1}{2})h) = y_{SM}^N((n-1)h)$$
 and $y_{SM}^N(0) = y(0) = y_0$.

The Strang-Marchuk method is a second-order method.

3. The Average Method

A new method is introduced here (referred to as the Average Method) based on the following idea: dividing the Cauchy problem (1) into d subproblems, using sequential splitting in all possible sequences, calculating the numerical solutions and then taking their arithmetic mean and letting it be the numerical solution in ω_h .

Let \mathcal{P}^n denote the set for the permutations of the indices $\{1, 2, \dots, n\}$. For $p = \{p_1, p_2, \dots, p_n\} \in \mathcal{P}^n$ we introduce the notation

$$\exp\{p_1, p_2, \dots, p_n\} = \exp(hA_{p_1}) \exp(hA_{p_2}) \cdot \dots \cdot \exp(hA_{p_n}). \tag{9}$$

Solving the Cauchy-problem (1) using sequential splitting for all possible permutations and then averaging the resulting numerical solutions yields a second-order method, i.e.

Statement 1.

$$\exp(h(A_1 + \dots + A_d)) = \frac{1}{d!} \sum_{p \in \mathcal{P}^d} \exp\{p_1, p_2, \dots, p_n\} + \mathcal{O}(h^3).$$
 (10)

See Appendix 1 for the proof.

Since \mathcal{P}^d has d! elements, we examine conditions that reduce the complexity of the Average Method. Consider the case of d=3, i.e. $A=A_1+A_2+A_3$. By Statement 1 we have

$$\exp(h(A_1 + \ldots + A_d)) = \frac{1}{3!} \sum_{p \in \mathcal{P}^3} \exp\{p_1, p_2, p_3\},\tag{11}$$

with 3! = 6 terms on the right-hand side. We utilize the usual definition of the commutator [A, B] = AB - BA. If, for example, A_1 and A_3 commute, i.e. $[A_1, A_3] = \mathbf{0}$, then we have

$$\exp\{p_2, p_1, p_3\} = \exp\{p_2, p_3, p_1\},\tag{12}$$

$$\exp\{p_1, p_3, p_2\} = \exp\{p_3, p_1, p_2\},\tag{13}$$

and

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$$\sum_{p \in \mathcal{P}^3} \exp\{p_1, p_2, p_3\} = \exp\{p_1, p_2, p_3\} + \exp\{p_3, p_2, p_1\} + 2\exp\{p_2, p_1, p_3\} + 2\exp\{p_1, p_3, p_2\}, \quad (14)$$

where the number of terms was reduced from six to four.

Let us now consider the general case. Let $A = A_1 + A_2 + ... + A_d$, and suppose that \exists $i, j \in \mathbb{N}, i \neq j$ such that $[A_i, A_j] = 0$. Then instead of all the d! permutations, we have d! - (d-1)! = (d-1)(d-1)! elements. If the decomposition includes more commuting pairs of matrices, the reduction might be more significant. An interesting question is the decomposition of a given matrix into pairwise commuting matrices. Even the study of the pairs of $n \times n$ commuting matrices A and B yields non-trivial results (i.e. Schur's theorem, Gerstenhaber's theorem, see, for example, Section 5 of [6]). The so-called commuting variety generated by the n^2 equations $(AB)_{ij} - (BA)_{ij} = 0$ can be investigated with the tools of algebraic geometry or linear algebraic conditions can be sought for commutativity. For k matrices one deals with commutative k-generated subalgebras, where even the best upper bound for their dimension is an open problem.

4. Making the Average Method third-order

Now we consider if third-order accuracy can be achieved with the Average Method based on sequential splitting. Assume that we have the Cauchy problem (1), with d=2. We then have the following

Statement 2. If and only if $A = A_1 + A_2$, and A_1 and A_2 satisfy the condition

$$[A_1, [A_1, A_2]] = [A_2, [A_1, A_2]]$$
 then

$$\exp(h(A_1 + A_2)) = \frac{\exp(hA_1)\exp(hA_2) + \exp(hA_2)\exp(hA_1)}{2} + \mathcal{O}(h^4).$$
 (15)

See Appendix 2 for the proof.

Now we consider the Strang-Marchuk splitting as the base method of average splitting. Can third-order accuracy be obtained by the Average Method?

The decomposition $A = A_1 + A_2$ with weights α and β $(\alpha, \beta \neq 0)$ and $\alpha + \beta \neq 0$ has to satisfy

$$\exp\left(h(A_1 + A_2)\right) =$$

$$= \frac{\alpha\left(\exp\left(\frac{hA_1}{2}\right)\exp(hA_2)\exp\left(\frac{hA_1}{2}\right)\right) + \beta\left(\exp\left(\frac{hA_2}{2}\right)\exp(hA_1)\exp\left(\frac{hA_2}{2}\right)\right)}{\alpha + \beta} + O(h^4). \quad (16)$$

The right-hand side of equation (16) can be written as

$$\frac{(\alpha+\beta)I + h(\alpha+\beta)(A_1 + A_2) + \frac{h^2}{2!}(\alpha+\beta)(A_1^2 + A_2^2 + A_1A_2 + A_2A_1)}{\alpha+\beta} + \frac{h^3}{3!} \left(\frac{(\alpha+\beta)(A_1^3 + A_2^3)}{\alpha+\beta} + \frac{(\frac{3}{4}\alpha + \frac{3}{2}\beta)A_1^2A_2 + (\frac{3}{4}\alpha + \frac{3}{2}\beta)A_2A_1^2 + (\frac{3}{2}\alpha + \frac{3}{4}\beta)A_2^2A_1}{\alpha+\beta} + \frac{(\frac{3}{2}\alpha + \frac{3}{4}\beta)A_1A_2^2 + \frac{3}{2}\alpha A_1A_2A_1 + \frac{3}{2}\beta A_2A_1A_2}{\alpha+\beta} \right) + O(h^4) =
= I + h(A_1 + A_2) + \frac{h^2}{2!}(A_1^2 + A_2^2 + A_1A_2 + A_2A_1) + \frac{h^3}{3!} \left((A_1^3 + A_2^3) + \frac{(\frac{3}{4}\alpha + \frac{3}{2}\beta)A_1^2A_2 + (\frac{3}{4}\alpha + \frac{3}{2}\beta)A_2A_1^2 + (\frac{3}{2}\alpha + \frac{3}{4}\beta)A_2^2A_1}{\alpha+\beta} + \frac{(\frac{3}{2}\alpha + \frac{3}{4}\beta)A_1A_2^2 + \frac{3}{2}\alpha A_1A_2A_1 + \frac{3}{2}\beta A_2A_1A_2}{\alpha+\beta} \right) + O(h^4).$$
(17)

The left-hand side of Equation (16) is

$$\exp\left(h(A_1 + A_2)\right) = I + h(A_1 + A_2) + \frac{h^2}{2!}(A_1 + A_2)^2 + \frac{h^3}{3!}(A_1 + A_2)^3 + \mathcal{O}(h^4) =$$

$$= I + h(A_1 + A_2) + \frac{h^2}{2}\left(A_1^2 + A_2^2 + A_1A_2 + A_2A_1\right) +$$

$$+ \frac{h^3}{6}\left(A_1^3 + A_2^3 + A_1A_2^2 + A_1^2A_2 + A_2A_1^2 + A_2^2A_1 + A_1A_2A_1 + A_2A_1A_2\right) + \mathcal{O}(h^4).$$
(18)

The equality (16) is only true under the conditions

$$\frac{\frac{3}{4}\alpha + \frac{3}{2}\beta}{\alpha + \beta} = 1,\tag{19}$$

$$\frac{\frac{3}{2}\alpha + \frac{3}{4}\beta}{\alpha + \beta} = 1,\tag{20}$$

$$\frac{\frac{3}{2}\alpha}{\alpha+\beta} = 1,\tag{21}$$

$$\frac{\frac{3}{2}\beta}{\alpha+\beta} = 1. \tag{22}$$

Equations (19) and (21) give the condition $\alpha = 2\beta$, while Equation (22) yields the condition $\beta = 2\alpha$. This implies $\alpha = \beta = 0$, which contradicts our assumption.

Third-order accuracy thus cannot be achieved with the Average Method based on Strang-Marchuk splitting.

5. Example Application

As seen in Section 3, the number of terms needed for the Average Method can be reduced by decomposing the underlying matrix into a set of matrices that have commuting elements.

We now investigate the efficacy of the three splitting methods discussed above on a physical problem. The model was chosen because of the structure of the matrices involved, i.e. sparse matrices whose decomposition into a partially commuting set was easy.

A piecewise-linear model of flutter was investigated in [7] and [8]. The affine model equations contain the three system matrices (k = 0, 1, 2)

$$A_k = \begin{pmatrix} 0 & 1 & 0 & 0 \\ -1 & -(p_1 + p_2 \mu c_k) & -\mu^2 c_k p_2 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & c_k \mu & -(p_4 - c_k \mu^2) & -p_3 \end{pmatrix},$$

with the model parameters given in Table 1 (see [7]) and $\mu \in (0, \infty)$ represents the nondimensional wind speed.

Table 1. Parameters of the model

Parameter	c_0	c_1	c_2	d_1	d_2	p_1	p_2	p_3	p_4
Value	5.932	-6.846	2.662	2.56	-0.2515	0.1485	0.0147	0.0540	0.2748

Motivated by this model, we consider the following 4-dimensional Cauchy problem

$$\begin{cases} \dot{\mathbf{x}}(t) = A_k \mathbf{x}(t), \\ \mathbf{x}(0) = \mathbf{x}_0. \end{cases}$$
 (23)

5.1. **Decompositions of matrix** A_k . All three A_k matrices have the same structure, thus we can discuss the symbolic decompositions of A_k . First, we analyze the decomposition

$$A_k = A_{k_{(1)}} + A_{k_{(2)}}, (24)$$

where

$$A_{k_{(1)}} = \begin{pmatrix} 0 & 1 & 0 & 0 \\ 0 & -(p_1 + p_2 \mu c_k) & -\mu^2 c_k p_2 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 0 & -p_3 \end{pmatrix}, \quad A_{k_{(2)}} = \begin{pmatrix} 0 & 0 & 0 & 0 \\ -1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & c_k \mu & -(p_4 - c_k \mu^2) & 0 \end{pmatrix}.$$

Clearly, $A_{k_{(1)}}$ is an upper triangular matrix. On the other hand, $A_{k_{(2)}}$ is a strictly lower triangular matrix and hence it is nilpotent. Therefore $(A_{k_{(2)}})^m = \mathbf{0}$ for m > 2. This means that the exponential of the matrix $A_{k_{(2)}}$ can be computed exactly. Thus, when we realize the splitting methods, the solution of the subproblem with this matrix can be calculated exactly. We can also define a decomposition in which the exponential of each matrix can be calculated exactly. For example, the decomposition

$$A_k = A_{k_{(1)}} + A_{k_{(2)}} + A_{k_{(3)}}, (25)$$

with

has this property because it is the sum of two nilpotent matrices and a diagonal matrix.

Now we define three decompositions which consist of commuting matrices. First, we consider the decomposition:

$$A_k = A_{k_{(1)}} + A_{k_{(2)}} + A_{k_{(3)}}, (26)$$

where

$$A_{k_{(3)}} = \begin{pmatrix} 0 & 0 & 0 & 0 \\ 0 & 0 & -\mu^2 c_k p_2 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & c_k \mu & 0 & 0 \end{pmatrix}.$$

 $A_{k_{(3)}}$ is a nilpotent matrix with $(A_{k_{(3)}})^m = \mathbf{0}$ for m > 2, hence its exponential can be given exactly. For the matrices $A_{k_{(1)}}$ and $A_{k_{(2)}}$ $A_{k_{(1)}}A_{k_{(2)}} = \mathbf{0}$ and $A_{k_{(2)}}A_{k_{(1)}} = \mathbf{0}$, so $[A_{k_{(1)}}, A_{k_{(2)}}] = \mathbf{0}$. On the one hand we win some adventage, due to the Statement 1, and on the other hand, this decomposition has a disadvantage, namely we lost the property of exact solvability.

In the following decomposition, the original matrix A_k is written as the sum of four matrices, two of them are commuting, and each has the property of exact solvability. The decomposition is the following:

$$A_k = A_{k_{(1)}} + A_{k_{(2)}} + A_{k_{(3)}} + A_{k_{(4)}}, (27)$$

where

$$A_{k_{(3)}} = \begin{pmatrix} 0 & 1 & 0 & 0 \\ 0 & 0 & -\mu^2 c_k p_2 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 0 & 0 \end{pmatrix}, \qquad A_{k_{(4)}} = \begin{pmatrix} 0 & 0 & 0 & 0 \\ -1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & c_k \mu & -(p_4 - c_k \mu^2) & 0 \end{pmatrix}.$$

 $A_{k_{(1)}}$ and $A_{k_{(2)}}$ are diagonal matrices (so they commute), therefore in both cases we can solve the subproblems exactly. Further, matrices $A_{k_{(3)}}$ and $A_{k_{(4)}}$ are nilpotent, so we can produce the exact solutions of the subproblems.

The last decomposition is the following:

$$A_k = A_{k_{(1)}} + A_{k_{(2)}} + A_{k_{(3)}} + A_{k_{(4)}}, (28)$$

where

In this set matrices $A_{k_{(1)}}$ and $A_{k_{(2)}}$ commute, and the exponential of $A_{k_{(1)}}$ can be computed exactly (it is actually a rotation matrix). Matrices $A_{k_{(3)}}$ and $A_{k_{(4)}}$ are nilpotent and diagonal matrices, already discussed in previous decompositions.

6. Numerical Experiments

For the numerical implementation of operator splitting, the question is the following: what kind of methods to use for computing the solutions of subproblems? Here we considered three possibilities.

The first and perhaps the most obvious choice is when at every step we solve the subproblems by a numerical method with the same order of accuracy as that of the splitting method. This means that for sequential splitting we have to use a first-order method, such as the explicit or implicit Euler method. For Strang-Marchuk splitting and the Average Method we have to use a second-order method, e.g. second-order Euler or trapezoidal method.

The solution of the subproblems can also be given by truncating the infinite exponential sequence at the order of the splitting method. Finally, we can produce exact solutions by using decompositions of matrix A_k to yield exactly solvable subproblems. We have already seen such decompositions in the previous section.

We performed the following eight experiments to solve the Cauchy problem (23):

E1: Sequential splitting with decomposition (24), solving the subproblems corresponding to $A_{k_{(1)}}$ and $A_{k_{(2)}}$ with explicit Euler method and exactly, respectively.

E2: Sequential splitting with decomposition (25), all subproblems solved exactly.

E3: Explicit Euler solution of the full problem (no splitting).

E4: Strang-Marchuk splitting with decomposition (24), solving the subproblems corresponding to $A_{k_{(1)}}$ and $A_{k_{(2)}}$ with improved Euler method and exactly, respectively.

E5: Strang-Marchuk splitting with decomposition (25), all subproblems solved exactly.

E6: Average Method with sequential splitting and decomposition (24), solving the subproblems corresponding to $A_{k_{(1)}}$ and $A_{k_{(2)}}$ with explicit Euler method and exactly, respectively.

E7: Average Method with sequential splitting and decomposition (26) all subproblems solved with explicit Euler method in parallel.

E8: Improved Euler solution of the full problem (no splitting).

We now detail the results of the above numerical experiments. The parameter value $\mu = 0.2$ and the initial condition $\mathbf{x}_0 = (1, 1, 1, 1)$ was used in all computations.

E1: Figure 1 shows component $x_3(t)$ of the solution on the time interval [0, 100]. The figure shows how the splitting solution (red line) approximates the exact solution (blue line) by reducing the step size h.

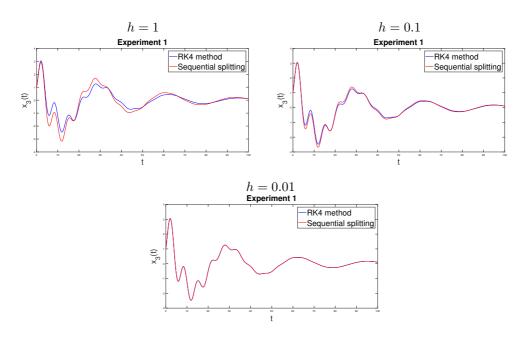


Figure 1. Experiment 1: $x_3(t)$ on time interval [0, 100]

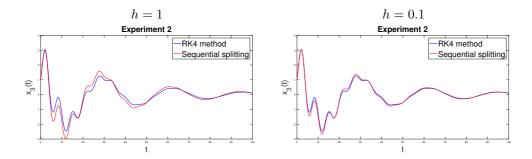


Figure 2. Experiment 2: $x_3(t)$ on the time interval [0, 100]

E2: In Figure 2 we see how the splitting solution approximates the exact solution by reducing the step size h. Figure 2 shows component x_3 of the solution on the time interval [0, 100]. The first 3 experiments utilized first-order methods. Runtimes are shown in Table 2.

Table 2. Comparison of runtimes (in seconds) for E1-3

h	1.	2.	3.	
	7.02×10^{-5}			
	8.44×10^{-4}			
0.01	1.70×10^{-3}	8.11×10^{-3}	1.09×10^{-2}	
0.001	1.37×10^{-2}	4.02×10^{-2}	7.14×10^{-1}	

As expected, we see that for the same order of accuracy splitting methods are faster than the full numerical solution. It can be seen that by reducing the step size h, the solvers containing splitting produce the numerical solution 1-2 orders of magnitude faster than the Euler method.

Table 3 shows the errors for Experiments 1 and 2. The methods give approximately the same error.

Table 3. Comparison of errors for Experiments 1 and 2

h	1.	2.
1.0	2.56×10^{-1}	2.64×10^{-1}
0.1	2.53×10^{-2}	2.09×10^{-2}
1		2.05×10^{-3}
0.001	2.08×10^{-4}	2.04×10^{-4}

E4: Since the Strang-Marchuk method is second-order, the choice of a bigger step size h is also sufficient to obtain a well-approximating splitting solution. This is illustrated in Figure 3, which shows $x_4(t)$ for the time interval [0, 100]. We see that even for h = 0.3 the splitting solution and the exact solution are very close.

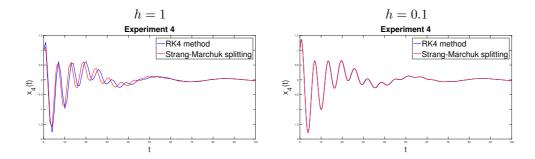


Figure 3. Experiment 4: $x_4(t)$ on the time interval [0, 100]

E5: We see in Figure 4 how the splitting solution behaves when the step size h is reduced. The figures show the component x_4 of the splitting solution on time interval [0, 100].

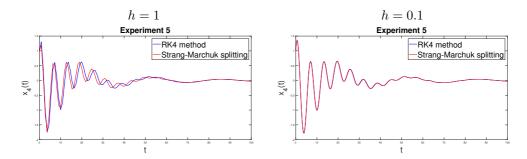


Figure 4. Experiment 5: $x_4(t)$ of the splitting solution on interval [0, 100]

For E4 and E5 the errors as well as the runtimes are similar. It is interesting to compare the errors between the Strang-Marchuk splitting method and using a second-order numerical method without splitting. We used the second-order Euler method to solve the system (23) without splitting in Experiment 8. In Table 4 we show the errors. The errors are two orders of magnitude smaller for the same stepsize h using the splitting method than using the second-order Euler method.

Table 4. Comparison of errors in case of Experiment 4 and 8

h	Experiment 4	Experiment 8
1.0	7.20×10^{-4}	8.21×10^{-2}
0.1	4.90×10^{-7}	8.24×10^{-5}
0.01	4.78×10^{-10}	7.52×10^{-8}
0.001	4.76×10^{-13}	7.43×10^{-11}

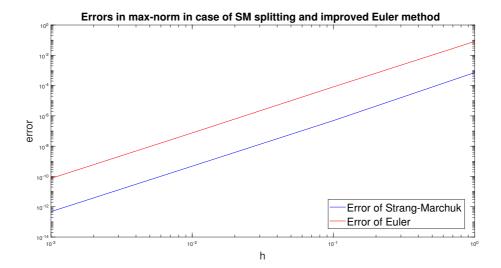


Figure 5. Error as a function of the stepsize h in Experiments 4 and 8

Figure 5 shows how errors change with decreasing h for Experiments 4 and 8.

E6: In this experiment we apply the average sequential splitting. At first we use the decomposition (24) requiring 2 numerical solutions. For this case, we have to solve the following two subproblems at each step:

$$\begin{cases} \dot{y_1}(t) = A_{0_{(1)}} y_1(t) & t \in [t_i, t_{i+1}], \\ y_1(t_i) = x_{sp(1)}(t_i), & (29) \end{cases} \begin{cases} \dot{y_2}(t) = A_{0_{(2)}} y_2(t) & t \in [t_i, t_{i+1}], \\ y_2(t_i) = y_1(t_{i+1}). & (30) \end{cases}$$

The splitting solution at t_{i+1} (i = 0, ..., n-1) is

$$x_{sp(1)}(t_{i+1}) = \exp\left(A_{0_{(2)}}(t_{i+1} - t_i)\right) y_1(t_{i+1}) \exp\left(A_{0_{(1)}}(t_{i+1} - t_i)\right) x_{sp(1)}(t_i).$$
(31)

We have to solve the following subproblems at each step:

$$\begin{cases}
\dot{y}_1(t) = A_{0(2)}y_1(t) & t \in [t_i, t_{i+1}], \\
y_1(t_i) = x_{sp(2)}(t_i), & (32)
\end{cases}
\begin{cases}
\dot{y}_2(t) = A_{0(1)}y_2(t) & t \in [t_i, t_{i+1}], \\
y_2(t_i) = y_1(t_{i+1}).
\end{cases}$$

The form of the splitting solution at t_{i+1} (i = 0, ..., n-1) is

$$x_{sp(2)}(t_{i+1}) = \exp(A_{0(1)}(t_{i+1} - t_i))y_1(t_{i+1}) \exp(A_{0(2)}(t_{i+1} - t_i))x_{sp(2)}(t_i).$$
(34)

The second-order approximate splitting solution is given by the average of the solutions $x_{sp(1)}(t_{i+1})$ and $x_{sp(2)}(t_{i+1})$, i.e.

$$x_{sp}(t_{i+1}) = \frac{x_{sp(1)}(t_{i+1}) + x_{sp(2)}(t_{i+1})}{2}.$$
(35)

Figure 6 shows how the splitting solution approximates the exact solution with decreasing h. The figures show $x_4(t)$ on the [0, 100] time interval.

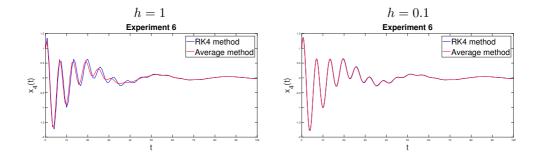


Figure 6. Experiment 6: $x_4(t)$ of the splitting solution on time interval [0, 100]

E7: We consider the decomposition (26) where the first two matrices are commuting. Therefore, for using the Average Method instead of 3! = 6 split problems we have to solve four subproblems, only.

The decomposition has the following form

$$A_0 = A_{0_{(3)}} + A_{0_{(4)}}, (36)$$

where we use the notation $A_{0_{(4)}}$ for the sum of the two commuting matrices, i.e. $A_{0_{(4)}}:=A_{0_{(1)}}+A_{0_{(2)}}$, which means that

$$A_{0_{(4)}} = \begin{pmatrix} 0 & 1 & 0 & 0 \\ -1 & -(p_1 + p_2 \mu c_k) & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & -(p_4 - c_k \mu^2) & -p_3 \end{pmatrix}.$$

Thus, in these cases, we solve the following sequences of sub-problems.

First, we solve the following two subproblems where the commutativity is not present.

The ordering $A_1 \to A_3 \to A_2$ results in the following split sub-problems:

$$\begin{cases} \dot{y_1}(t) = A_{0_{(1)}} y_1(t) & t \in [t_i, t_{i+1}], \\ y_1(t_i) = x_{sp(1)}(t_i), & \\ \end{cases} \begin{cases} \dot{y_2}(t) = A_{0_{(3)}} y_2(t) & t \in [t_i, t_{i+1}], \\ y_2(t_i) = y_1(t_{i+1}), & \\ \end{cases}$$
(38)

$$\begin{cases}
\dot{y}_3(t) = A_{0_{(2)}} y_3(t) & t \in [t_i, t_{i+1}], \\
y_3(t_i) = y_2(t_{i+1}).
\end{cases}$$
(39)

Then the splitting solution at $t = t_{i+1}$ is

$$x_{sp(1)}(t_{i+1}) = \exp\left(A_{0_{(2)}}(t_{i+1} - t_i)\right) y_2(t_{i+1}) \exp\left(A_{0_{(3)}}(t_{i+1} - t_i)\right) y_1(t_{i+1}) \exp\left(A_{0_{(1)}}(t_{i+1} - t_i)\right) x_{sp(1)}(t_i).$$

$$(40)$$

The ordering $A_2 \to A_3 \to A_1$ results in the sub-problems following three subproblems:

$$\begin{cases} \dot{y_1}(t) = A_{0_{(2)}} y_1(t) & t \in [t_i, t_{i+1}], \\ y_1(t_i) = x_{sp(2)}(t_i), & \\ \end{cases} \begin{cases} \dot{y_2}(t) = A_{0_{(3)}} y_2(t) & t \in [t_i, t_{i+1}], \\ y_2(t_i) = y_1(t_{i+1}), & \\ \end{cases}$$

$$(42)$$

$$\begin{cases}
\dot{y_3}(t) = A_{0_{(1)}} y_1(t) & t \in [t_i, t_{i+1}], \\
y_3(t_i) = y_2(t_{i+1}).
\end{cases}$$
(43)

Then the splitting solution at t_{i+1} is

$$x_{sp(2)}(t_{i+1}) = \exp \left(A_{0_{(1)}}(t_{i+1} - t_i)\right) y_2(t_{i+1}) \exp \left(A_{0_{(3)}}(t_{i+1} - t_i)\right) y_1(t_{i+1}) \exp \left(A_{0_{(2)}}(t_{i+1} - t_i)\right) x_{sp(2)}(t_i).$$

$$(44)$$

Due to the commutativity, the remaining two problems, which we solve, consist of only two subproblems. These are the following.

For the ordering $A_4 \to A_3$ the subproblems are the following:

$$\begin{cases} \dot{y_1}(t) = A_{0_{(4)}} y_1(t) & t \in [t_i, t_{i+1}], \\ y_1(t_i) = x_{sp(3)}(t_i), & (45) \end{cases} \begin{cases} \dot{y_2}(t) = A_{0_{(3)}} y_2(t) & t \in [t_i, t_{i+1}], \\ y_2(t_i) = y_1(t_{i+1}). & (46) \end{cases}$$

The splitting solution at t_{i+1} is

$$x_{sp(3)}(t_{i+1}) = \\ = \exp\left(A_{0_{(3)}}(t_{i+1} - t_i)\right) y_1(t_{i+1}) \exp\left(A_{0_{(4)}}(t_{i+1} - t_i)\right) x_{sp(3)}(t_i). \tag{47}$$
 Finally, for the ordering $A_3 \to A_4$ we get

$$\begin{cases} \dot{y_1}(t) = A_{0_{(3)}} y_1(t) & t \in [t_i, t_{i+1}], \\ y_1(t_i) = x_{sp(4)}(t_i), & (48) \end{cases} \begin{cases} \dot{y_2}(t) = A_{0_{(4)}} y_2(t) & t \in [t_i, t_{i+1}], \\ y_2(t_i) = y_1(t_{i+1}). \end{cases}$$

Then

$$x_{sp(4)}(t_{i+1}) =$$

$$= \exp\left(A_{0_{(4)}}(t_{i+1} - t_i)\right) y_1(t_{i+1}) \exp\left(A_{0_{(3)}}(t_{i+1} - t_i)\right) x_{sp(4)}(t_i). \quad (50)$$

Using the Average Method, the second-order accurate approximation is defined as

$$x_{sp}(t) = \frac{x_{sp(1)}(t) + x_{sp(2)}(t) + 2 \cdot x_{sp(3)}(t) + 2 \cdot x_{sp(4)}(t)}{6}.$$
 (51)

Figure 7 shows how the 4th component $(x_4(t))$ of the splitting solution approaches the exact solution with decreasing h. The advantage of the method is

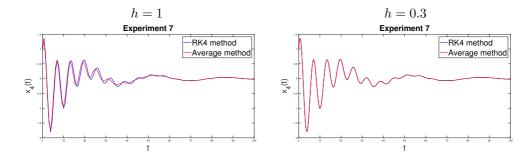


Figure 7. Experiment 7: $x_4(t)$ on time interval interval [0, 100]

that the solutions (40), (44), (47) and (50) can be independently calculated, i.e., the computation is parallelizable.

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table 5.	Comparison of	oi runtimes	(in seconds) for Experiments 4-8.

h	E 4	E 5	E 6	E7	E8
1.0		7.57×10^{-2}			
0.1	7.55×10^{-1}	8.71×10^{-1}	1.01×10^{-3}	1.52×10^{-3}	1.96×10^{-2}
0.01	5.20×10^{0}	6.39×10^{0}	3.65×10^{-3}	7.81×10^{-3}	8.44×10^{-2}
0.001	1.53×10^{1}	2.06×10^{1}	1.89×10^{-2}	6.47×10^{-2}	1.13×10^{0}

Table 5 collects the comparison of runtimes for Experiments 4-8. These results show that sequential and Strang-Marchuk splitting is about two orders of magnitude slower than the improved Euler method, while the Average Method is about two orders of magnitude faster. This is an encouraging result for the applicability of the Average Method.

7. Summary

By performing several numerical experiments we demonstrated that the benefits of the Average Method are the following:

- easy implementation when d is small,
- provides a second-order approximation solution using a first-order method,
- the numerical solutions of the subproblems can be independently computed, therefore the method can be parallelized.

The drawback of the Average Method is the large number (d!) of solutions to be computed. This number can be reduced by finding a decomposition of the system matrix into a set of pairwise commuting matrices.

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APPENDIX 1

Statement 1. Assume that we have the Cauchy problem (1) Solving this problem in all possible sequences using sequential splitting, and then taking the average of the resulting numerical solutions, the method is second-order, i.e.

$$\exp(h(A_1 + \dots + A_d)) = \frac{1}{d!} \sum_{p \in \mathcal{P}^d} \exp\{p_1, p_2, \dots, p_n\} + \mathcal{O}(h^3).$$
 (10)

Proof. We prove the statement by induction. First we consider the case k=2, i.e. the validity of the formula

$$\exp(h(A_1 + A_2)) = \frac{\exp(hA_1)\exp(hA_2) + \exp(hA_2)\exp(hA_1)}{2!} + \mathcal{O}(h^3).$$
 (52)

Obviously, for any matrix C we have

$$\exp(hC) = I + hC + \frac{h^2}{2}C^2 + \mathcal{O}(h^3).$$
 (53)

Hence,

$$\frac{\exp(hA_1) \exp(hA_2) + \exp(hA_2) \exp(hA_1)}{2!} = \frac{\left[I + hA_1 + \frac{h^2}{2!}A_1^2\right] \left[I + hA_2 + \frac{h^2}{2!}A_2^2\right] + \left[I + hA_2 + \frac{h^2}{2!}A_2^2\right] \left[I + hA_1 + \frac{h^2}{2!}A_1^2\right]}{2!} = \frac{2I + 2h(A_1 + A_2) + h^2(A_1^2 + A_2^2 + A_1A_2 + A_2A_1)}{2!} = \frac{2I + h(A_1 + A_2) + h^2(A_1^2 + A_2^2 + A_1A_2 + A_2A_1)}{2!} = \frac{I + h(A_1 + A_2) + \frac{h^2}{2!}(A_1^2 + A_2^2 + A_1A_2 + A_2A_1)}{2!} = \frac{I + h(A_1 + A_2) + \frac{h^2}{2!}(A_1^2 + A_2^2 + A_1A_2 + A_2A_1)}{2!} = \frac{I + h(A_1 + A_2) + \frac{h^2}{2!}(A_1^2 + A_2^2 + A_1A_2 + A_2A_1)}{2!} = \frac{I + h(A_1 + A_2) + \frac{h^2}{2!}(A_1^2 + A_2^2 + A_1A_2 + A_2A_1)}{2!} = \frac{I + h(A_1 + A_2) + \frac{h^2}{2!}(A_1^2 + A_2^2 + A_1A_2 + A_2A_1)}{2!} = \frac{I + h(A_1 + A_2) + \frac{h^2}{2!}(A_1^2 + A_2^2 + A_1A_2 + A_2A_1)}{2!} = \frac{I + h(A_1 + A_2) + \frac{h^2}{2!}(A_1^2 + A_2^2 + A_1A_2 + A_2A_1)}{2!} = \frac{I + h(A_1 + A_2) + \frac{h^2}{2!}(A_1^2 + A_2^2 + A_1A_2 + A_2A_1)}{2!} = \frac{I + h(A_1 + A_2) + \frac{h^2}{2!}(A_1^2 + A_2^2 + A_1A_2 + A_2A_1)}{2!} = \frac{I + h(A_1 + A_2) + \frac{h^2}{2!}(A_1^2 + A_2^2 + A_1A_2 + A_2A_1)}{2!} = \frac{I + h(A_1 + A_2) + \frac{h^2}{2!}(A_1^2 + A_2^2 + A_1A_2 + A_2A_1)}{2!} = \frac{I + h(A_1 + A_2) + \frac{h^2}{2!}(A_1^2 + A_2^2 + A_1A_2 + A_2A_1)}{2!} = \frac{I + h(A_1 + A_2) + \frac{h^2}{2!}(A_1^2 + A_2^2 + A_1A_2 + A_2A_1)}{2!} = \frac{I + h(A_1 + A_2) + \frac{h^2}{2!}(A_1^2 + A_2^2 + A_1A_2 + A_2A_1)}{2!} = \frac{I + h(A_1 + A_2) + \frac{h^2}{2!}(A_1^2 + A_2^2 + A_1A_2 + A_2A_1)}{2!} = \frac{I + h(A_1 + A_2) + \frac{h^2}{2!}(A_1^2 + A_2^2 + A_1A_2 + A_2A_1)}{2!} = \frac{I + h(A_1 + A_2) + \frac{h^2}{2!}(A_1^2 + A_2^2 + A_1A_2 + A_2A_1)}{2!} = \frac{I + h(A_1 + A_2) + \frac{h^2}{2!}(A_1^2 + A_2^2 + A_1A_2 + A_2A_1)}{2!} = \frac{I + h(A_1 + A_2) + \frac{h^2}{2!}(A_1^2 + A_2^2 + A_1A_2 + A_2A_1)}{2!} = \frac{I + h(A_1 + A_2) + \frac{h^2}{2!}(A_1^2 + A_2^2 + A_1A_2 + A_2A_1)}{2!} = \frac{I + h(A_1 + A_2) + \frac{h^2}{2!}(A_1^2 + A_2^2 + A_1A_2 + A_2A_1)}{2!} = \frac{I + h(A_1 + A_2) + \frac{h^2}{2!}(A_1^2 + A_2^2 + A_1A_2 + A_2A_1)}{2!} = \frac{I + h(A_1 + A_2) + \frac{h^2}{2!}(A_1^2 + A_2^2 + A_1A_2 + A_2A_1)}{2!} = \frac{I + h(A_1 + A_2) + \frac{h^2}{2!}(A_1^2 + A_2^2 + A_1A_2 + A_2A_1)}{2!} = \frac{I$$

which proves (52).

Now we suppose that the statement is true for k = d and prove its validity for k = d + 1.

$$\exp\left(h(A_1 + A_2 + \dots + A_d + A_{d+1})\right) =$$

$$= I + h \sum_{j=1}^{d+1} A_j + \frac{h^2}{2!} \left(\sum_{j=1}^{d+1} A_j^2 + \sum_{i=1}^{d} \sum_{j=i+1}^{d+1} A_i A_j + \sum_{i=2}^{d+1} \sum_{j=1}^{i-1} A_i A_j\right) + \mathcal{O}(h^3). \quad (55)$$

Using the notation $B = A_1 + A_2 + \ldots + A_d$, we have

$$\exp\left(h(A_{1} + A_{2} + \dots + A_{d} + A_{d+1})\right) =$$

$$= \exp\left(h(B + A_{d+1})\right) = \frac{\exp(hB) \exp(hA_{d+1}) + \exp(hA_{d+1}) \exp(hB)}{2!} + \mathcal{O}(h^{3}) =$$

$$= \frac{1}{2} \left[I + h \sum_{j=1}^{d} A_{j} + \frac{h^{2}}{2!} \left(\sum_{j=1}^{d} A_{j}^{2} + \sum_{i=1}^{d-1} \sum_{j=i+1}^{d} A_{i} A_{j} + \sum_{i=2}^{d} \sum_{j=1}^{i-1} A_{i} A_{j}\right) +$$

$$+ hA_{d+1} + h^{2} \sum_{j=1}^{d} A_{j} A_{d+1} + \frac{h^{2}}{2} A_{d+1}^{2}\right] +$$

$$+ \frac{1}{2} \left[I + h \sum_{j=1}^{d} A_{j} + \frac{h^{2}}{2!} \left(\sum_{j=1}^{d} A_{j}^{2} + \sum_{i=1}^{d-1} \sum_{j=i+1}^{d} A_{i} A_{j} + \sum_{i=2}^{d} \sum_{j=1}^{i-1} A_{i} A_{j}\right) +$$

$$+ hA_{d+1} + h^{2} \sum_{j=1}^{d} A_{d+1} A_{j} + \frac{h^{2}}{2} A_{d+1}^{2}\right] + \mathcal{O}(h^{3}) =$$

$$= \frac{1}{2} \left[2I + 2h \left(\sum_{j=1}^{d} A_{j} + A_{d+1}\right) + h^{2} \left(\sum_{j=1}^{d} A_{j}^{2} + A_{d+1}^{2} + \sum_{i=1}^{d-1} \sum_{j=i+1}^{d} A_{i} A_{j} + \sum_{i=1}^{d-1} \sum_{j=i+1}^{d} A_{i} A_{j} + \sum_{i=1}^{d+1} \sum_{j=1}^{d+1} A_{i} A_{j} + \sum_{i=2}^{d+1} \sum_{j=1}^{d+1} A_{i} A_{j} + \sum_{i=2}^{d+1} \sum_{j=1}^{d+1} A_{i} A_{j} + \sum_{i=2}^{d+1} \sum_{j=1}^{d+1} A_{i} A_{j}\right] + \mathcal{O}(h^{3}) =$$

$$= I + h \sum_{j=1}^{d+1} A_{j} + \frac{h^{2}}{2} \left(\sum_{j=1}^{d+1} A_{j}^{2} + \sum_{i=1}^{d} \sum_{j=i+1}^{d+1} A_{i} A_{j} + \sum_{i=2}^{d+1} \sum_{j=1}^{d+1} A_{i} A_{j} + \sum_{i=1}^{d+1} A_{i} A_{$$

This proves our statement. \square

Appendix 2

Statement 2. If and only if $A = A_1 + A_2$, and A_1 and A_2 satisfy the condition $\begin{bmatrix} A_1, [A_1, A_2] \end{bmatrix} = \begin{bmatrix} A_2, [A_1, A_2] \end{bmatrix}$ then

$$\exp(h(A_1 + A_2)) = \frac{\exp(hA_1)\exp(hA_2) + \exp(hA_2)\exp(hA_1)}{2} + \mathcal{O}(h^4).$$
 (15)

Proof. Let $x, y \in \mathbb{R} \setminus \{0\}$, $x + y \neq 0$ and d = 2 and we want to prove the following relation:

$$\exp(h(A_1 + A_2)) = \frac{x \exp(hA_1) \exp(hA_2) + y \exp(hA_2) \exp(hA_1)}{x + y} + \mathcal{O}(h^4).$$
 (57)

Clearly

$$\exp\left(h(A_1 + A_2)\right) =$$

$$= I + h(A_1 + A_2) + \frac{h^2}{2!}(A_1 + A_2)^2 + \frac{h^3}{3!}(A_1 + A_2)^3 + \mathcal{O}(h^4) =$$

$$= I + h(A_1 + A_2) + \frac{h^2}{2}\left(A_1^2 + A_2^2 + A_1A_2 + A_2A_1\right) + \frac{h^3}{6}\left(A_1^3 + A_2^3 + A_1A_2^2 + A_1A_2 + A_2A_1^2 + A_2A_1 + A_2A_1A_2\right) + \mathcal{O}(h^4). \quad (58)$$

Then, the right-hand side of (57):

$$\frac{x \exp(hA_1) \exp(hA_2) + y \exp(hA_2) \exp(hA_1)}{x + y} + \mathcal{O}(h^4) = \frac{1}{x + y} \left[x \left(I + hA_1 + \frac{h^2}{2!} A_1^2 + \frac{h^3}{3!} A_1^3 + \mathcal{O}(h^4) \right) \left(I + hA_2 + \frac{h^2}{2!} A_2^2 + \frac{h^3}{3!} A_2^3 + \mathcal{O}(h^4) \right) \right] + \frac{1}{x + y} \left[y \left(I + hA_2 + \frac{h^2}{2!} A_2^2 + \frac{h^3}{3!} A_3^3 + \mathcal{O}(h^4) \right) \left(I + hA_1 + \frac{h^2}{2!} A_1^2 + \frac{h^3}{3!} A_1^3 + \mathcal{O}(h^4) \right) \right] = \frac{1}{x + y} \left[x \left(I + h(A_1 + A_2) + h^2 \left(\frac{1}{2} A_1^2 + \frac{1}{2} A_2^2 + A_1 A_2 \right) + h^3 \left(\frac{1}{6} A_1^3 + \frac{1}{6} A_2^3 + \frac{1}{2} A_1^2 A_2 + \frac{1}{2} A_1 A_2^2 \right) \right) \right] + \frac{1}{x + y} \left[y \left(I + h(A_1 + A_2) + h^2 \left(\frac{1}{2} A_1^2 + \frac{1}{2} A_2^2 + A_2 A_1 \right) + h^3 \left(\frac{1}{6} A_1^3 + \frac{1}{6} A_2^3 + \frac{1}{2} A_2^2 A_1 + \frac{1}{2} A_2 A_1^2 \right) \right) \right] = 1 + h(A_1 + A_2) + \frac{1}{x + y} \left\{ h^2 \left[x \left(\frac{1}{2} A_1^2 + \frac{1}{2} A_2^2 + A_1 A_2 \right) + y \left(\frac{1}{2} A_1^2 + \frac{1}{2} A_2^2 + A_2 A_1 \right) \right] + h^3 \left[x \left(\frac{1}{6} A_1^3 + \frac{1}{6} A_2^3 + \frac{1}{2} A_1^2 A_2 + \frac{1}{2} A_1 A_2^2 \right) + y \left(\frac{1}{6} A_1^3 + \frac{1}{6} A_2^3 + \frac{1}{2} A_2^2 A_1 + \frac{1}{2} A_2 A_1^2 \right) \right] \right\}. \tag{59}$$

(57) will be true if coefficients of (58) and (59) are the same. The coefficients of h^i , i = 0, 1, 2, 3 are the following.

Coefficient of h^0 :

$$\frac{(x+y)I}{(x+y)} = I \Longrightarrow I = I.$$

Coefficient of h^1 :

$$\frac{(x+y)(A_1+A_2)}{(x+y)} = (A_1+A_2) \Longrightarrow (A_1+A_2) = (A_1+A_2).$$

Coefficient of h^2 :

$$x\left(\frac{1}{2}A_{1}^{2} + \frac{1}{2}A_{2}^{2} + A_{1}A_{2}\right) + y\left(\frac{1}{2}A_{1}^{2} + \frac{1}{2}A_{2}^{2} + A_{2}A_{1}\right) =$$

$$= \frac{1}{2}(A_{1} + A_{2}) + \frac{1}{2}(A_{1}A_{2} + A_{2}A_{1})$$

$$\frac{1}{2}(A_{1} + A_{2}) + \frac{xA_{1}A_{2} + yA_{2}A_{1}}{x + y} = \frac{1}{2}(A_{1} + A_{2}) + \frac{1}{2}(A_{1}A_{2} + A_{2}A_{1}),$$

$$\frac{xA_{1}A_{2} + yA_{2}A_{1}}{x + y} = \frac{1}{2}(A_{1}A_{2} + A_{2}A_{1}),$$

$$xA_{1}A_{2} + yA_{2}A_{1} = \frac{1}{2}(x + y)A_{1}A_{2} + \frac{1}{2}(x + y)A_{2}A_{1}.$$

$$\downarrow \qquad \qquad \downarrow \qquad \downarrow \qquad \qquad$$

Coefficient of h^3 :

$$\frac{1}{x+y} \left[x \left(\frac{1}{6} A_1^3 + \frac{1}{6} A_2^3 + \frac{1}{2} A_1^2 A_2 + \frac{1}{2} A_1 A_2^2 \right) + y \left(\frac{1}{6} A_1^3 + \frac{1}{6} A_2^3 + \frac{1}{2} A_2^2 A_1 + \frac{1}{2} A_2 A_1^2 \right) \right] =
= \frac{1}{6} \left(A_1^3 + A_2^3 + A_1 A_2^2 + A_1^2 A_2 + A_2 A_1^2 + A_2^2 A_1 + A_1 A_2 A_1 + A_2 A_1 A_2 \right).$$

After some calculations we get

$$\frac{x(A_1^2A_2 + A_1A_2^2) + y(A_2^2A_1 + A_2A_1^2)}{2(x+y)} = \frac{A_1A_2^2 + A_1^2A_2 + A_2A_1^2 + A_2^2A_1 + A_1A_2A_1 + A_2A_1A_2}{6}$$

Utilizing equation (60) yields

$$\frac{x(A_1^2A_2 + A_1A_2^2 + A_2^2A_1 + A_2A_1^2)}{4x} = \frac{A_1A_2^2 + A_1^2A_2 + A_2A_1^2 + A_2^2A_1 + A_1A_2A_1 + A_2A_1A_2}{6}$$

or

$$A_1^2 A_2 + A_1 A_2^2 + A_2^2 A_1 + A_2 A_1^2 = 2A_1 A_2 A_1 + 2A_2 A_1 A_2$$

and

$$A_1[A_1, A_2] + [A_2, A_1]A_1 + A_2[A_2, A_1] + [A_1, A_2]A_2 = 0.$$

On the basis of the well known equality [A, B] = -[B, A] we have

$$A_1[A_1, A_2] - [A_1, A_2]A_1 - A_2[A_1, A_2] + [A_1, A_2]A_2 = 0,$$

or

$$A_1, [A_1, A_2] + [A_1, A_2], A_2 = 0,$$

and

$$\left[A_1, \, \left[A_1, A_2 \right] \right] - \left[A_2, \, \left[A_1, \, A_2 \right] \right] = 0.$$

Hence, the condition of third order is as follows

$$\[A_1, [A_1, A_2]\] = \[A_2, [A_1, A_2]\]. \tag{61}$$

This means that for a decomposition of two matrices the method will be third-order if and only if condition (61) is satisfied. \Box

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